Genuine Bell nonlocality in many-body quantum systems

Bachelor thesis written under supervision of dr hab. inż. Remigiusz Augusiak and dr hab. Rafał Demkowicz-Dobrzański

Ignacy Stachura

◆□▶ ◆□▶ ◆三▶ ◆三▶ ○ ● ●

Genuine multipartite entanglement

$$\rho_{AB} = \sum_{i} p_{i} \left| \psi_{A}^{i} \right\rangle \! \left\langle \psi_{A}^{i} \right| \otimes \left| \phi_{B}^{i} \right\rangle \! \left\langle \phi_{B}^{i} \right|, \quad p_{i} \ge 0, \quad \sum_{i} p_{i} = 1$$

• We consider a system $A := A^1, \ldots, A^N$ consisting of N subsystems.

$$\rho_{AB} = \sum_{i} p_{i} \left| \psi_{A}^{i} \right\rangle \! \left\langle \psi_{A}^{i} \right| \otimes \left| \phi_{B}^{i} \right\rangle \! \left\langle \phi_{B}^{i} \right|, \quad p_{i} \ge 0, \quad \sum_{i} p_{i} = 1$$

• We consider a system $A := A^1, \ldots, A^N$ consisting of N subsystems.

■ We divide the set of subsystems into two non-empty subsets: a subset *g* and its complement denoted as *g*.

$$\rho_{AB} = \sum_{i} p_{i} \left| \psi_{A}^{i} \right\rangle \! \left\langle \psi_{A}^{i} \right| \otimes \left| \phi_{B}^{i} \right\rangle \! \left\langle \phi_{B}^{i} \right|, \quad p_{i} \ge 0, \quad \sum_{i} p_{i} = 1$$

- We consider a system $A := A^1, \ldots, A^N$ consisting of N subsystems.
- We divide the set of subsystems into two non-empty subsets: a subset *g* and its complement denoted as *g*.
- We ask whether the state is separable across a given partition $g|\bar{g}$.

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

$$\rho_{AB} = \sum_{i} p_{i} \left| \psi_{A}^{i} \right\rangle \! \left\langle \psi_{A}^{i} \right| \otimes \left| \phi_{B}^{i} \right\rangle \! \left\langle \phi_{B}^{i} \right|, \quad p_{i} \ge 0, \quad \sum_{i} p_{i} = 1$$

- We consider a system $A := A^1, \ldots, A^N$ consisting of N subsystems.
- We divide the set of subsystems into two non-empty subsets: a subset *g* and its complement denoted as *g*.
- We ask whether the state is separable across a given partition $g|\bar{g}$.

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三 のへぐ

• We consider all the possible partitions.

$$\rho_{AB} = \sum_{i} p_{i} \left| \psi_{A}^{i} \right\rangle \! \left\langle \psi_{A}^{i} \right| \otimes \left| \phi_{B}^{i} \right\rangle \! \left\langle \phi_{B}^{i} \right|, \quad p_{i} \ge 0, \quad \sum_{i} p_{i} = 1$$

- We consider a system $A := A^1, \ldots, A^N$ consisting of N subsystems.
- We divide the set of subsystems into two non-empty subsets: a subset g and its complement denoted as ḡ.
- We ask whether the state is separable across a given partition $g|\bar{g}$.
- We consider all the possible partitions.

$$egin{aligned}
ho_\mathsf{A} &= \sum_{g \mid ar{g}} q_{g \mid ar{g}} \sum_i p_i^{g \mid ar{g}} \left| \psi_g^i
ight
angle \psi_g^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \ , \ & \sum_{g \mid ar{g}} q_{g \mid ar{g}} = 1, \quad \forall_{g \mid ar{g}} \sum_i p_i^{g \mid ar{g}} = 1 \end{aligned}$$

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三 のへぐ

$$\rho_{AB} = \sum_{i} p_{i} \left| \psi_{A}^{i} \right\rangle \! \left\langle \psi_{A}^{i} \right| \otimes \left| \phi_{B}^{i} \right\rangle \! \left\langle \phi_{B}^{i} \right|, \quad p_{i} \ge 0, \quad \sum_{i} p_{i} = 1$$

- We consider a system $A := A^1, \ldots, A^N$ consisting of N subsystems.
- We divide the set of subsystems into two non-empty subsets: a subset g and its complement denoted as ḡ.
- We ask whether the state is separable across a given partition $g|\bar{g}$.
- We consider all the possible partitions.

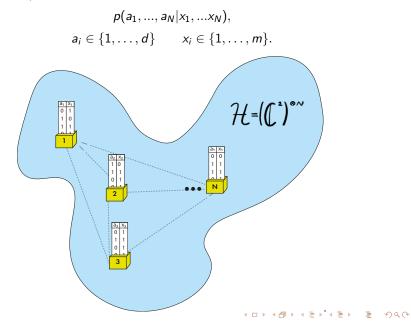
$$egin{aligned}
ho_\mathsf{A} &= \sum_{g \mid ar{g}} q_{g \mid ar{g}} \sum_i p_i^{g \mid ar{g}} \left| \psi_g^i
ight
angle \psi_g^i
ight
angle \psi_g^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i \left| \otimes \left| \psi_{ar{g}}^i
ight
angle \psi_{ar{g}}^i
ight$$

GME

States that do not admit the above are GME states.

Bell-like experiment

In (N, m, d) scenario we obtain the conditional probabilities



Genuine multipartite nonlocality

$$p(a_1, a_2|x_1, x_2) = \sum_{\lambda} q(\lambda) p(a_1|x_1, \lambda) p(a_2|x_2, \lambda)$$

• Correlations bilocal across a partition $g|\bar{g}$:

$$p(\vec{a}|\vec{x}) = \sum_{\lambda} q(\lambda) p(\vec{a}_g|\vec{x}_{\bar{g}},\lambda) p(\vec{a}_{\bar{g}}|\vec{x}_{\bar{g}},\lambda).$$

(ロト (個) (E) (E) (E) (E) のへの

Genuine multipartite nonlocality

$$p(a_1,a_2|x_1,x_2) = \sum_{\lambda} q(\lambda) p(a_1|x_1,\lambda) p(a_2|x_2,\lambda)$$

• Correlations bilocal across a partition $g|\bar{g}$:

$$p(\vec{a}|\vec{x}) = \sum_{\lambda} q(\lambda) p(\vec{a}_{g}|\vec{x}_{\bar{g}},\lambda) p(\vec{a}_{\bar{g}}|\vec{x}_{\bar{g}},\lambda).$$

Just bilocal correlations:

$$egin{aligned} p(ec{a}ec{x}) &= \sum_{gec{g}}\sum_{\lambda} q_g(\lambda) p(ec{a}_gec{x}_g,\lambda) p(ec{a}_{ec{g}}ec{x}_{ec{g}},\lambda), \ &\sum \sum q_g(\lambda) = 1. \end{aligned}$$

▲□▶ ▲□▶ ▲ 臣▶ ▲ 臣▶ ― 臣 … 釣��

 $g|\bar{g} \lambda$

Genuine multipartite nonlocality

$$p(a_1, a_2|x_1, x_2) = \sum_{\lambda} q(\lambda) p(a_1|x_1, \lambda) p(a_2|x_2, \lambda)$$

• Correlations bilocal across a partition $g|\bar{g}$:

$$p(\vec{a}|\vec{x}) = \sum_{\lambda} q(\lambda) p(\vec{a}_{g}|\vec{x}_{\bar{g}},\lambda) p(\vec{a}_{\bar{g}}|\vec{x}_{\bar{g}},\lambda).$$

Just bilocal correlations:

$$egin{aligned} p(ec{a}ec{ec{x}}) &= \sum_{gec{g}}\sum_{\lambda}q_g(\lambda)p(ec{a}_gec{ec{x}}_g,\lambda)p(ec{a}_{ec{g}}ec{ec{x}}_{ec{g}},\lambda), \ &\sum_{gec{g}}\sum_{\lambda}q_g(\lambda) = 1. \end{aligned}$$

GMNL

Correlations that are not bilocal are GMNL.

More subsystems

▲□▶ ▲□▶ ▲ 臣▶ ▲ 臣▶ ― 臣 … 釣��

 There exist mixed 2-qudit entangled states that are local.

More subsystems

▲□▶ ▲□▶ ▲ 臣▶ ▲ 臣▶ ― 臣 … のへぐ

- There exist mixed 2-qudit entangled states that are local.
- All pure 2-qudit entangled states are nonlocal.

- There exist mixed 2-qudit entangled states that are local.
- All pure 2-qudit entangled states are nonlocal.

More subsystems

• There exist mixed GME states that are not GMNL.

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三 のへぐ

- There exist mixed 2-qudit entangled states that are local.
- All pure 2-qudit entangled states are nonlocal.

More subsystems

• There exist mixed GME states that are not GMNL.

◆□▶ ◆□▶ ◆三▶ ◆三▶ ○ ● ●

• Conjecture: All pure GME states are GMNL.

- There exist mixed 2-qudit entangled states that are local.
- All pure 2-qudit entangled states are nonlocal.

More subsystems

- There exist mixed GME states that are not GMNL.
- Conjecture: All pure GME states are GMNL.
 - 3-qubit states with symmetry between 2 qubits

▲ロ ▶ ▲周 ▶ ▲ 国 ▶ ▲ 国 ▶ ● の Q @

- There exist mixed 2-qudit entangled states that are local.
- All pure 2-qudit entangled states are nonlocal.

More subsystems

- There exist mixed GME states that are not GMNL.
- Conjecture: All pure GME states are GMNL.
 - 3-qubit states with symmetry between 2 qubits
 - special cases e.g. GHZ states for any n

▲ロ ▶ ▲周 ▶ ▲ 国 ▶ ▲ 国 ▶ ● の Q @

- There exist mixed 2-qudit entangled states that are local.
- All pure 2-qudit entangled states are nonlocal.

More subsystems

- There exist mixed GME states that are not GMNL.
- Conjecture: All pure GME states are GMNL.
 - 3-qubit states with symmetry between 2 qubits
 - special cases e.g. GHZ states for any n

▲ロ ▶ ▲周 ▶ ▲ 国 ▶ ▲ 国 ▶ ● の Q @

 some numerical results e.g 3-qubit states F. J. Curchod et al., New J. Phys. 21, 023016 (2019)

CHSH inequality

 $IAB = p(00|00) - p(01|01) - p(10|10) - p(00|11) \le 0$

◆□▶ ◆□▶ ◆三▶ ◆三▶ ○ ● ●

F. J. Curchod et al., New J. Phys. 21, 023016 (2019)

CHSH inequality

$$IAB = p(00|00) - p(01|01) - p(10|10) - p(00|11) \le 0$$

Lifted term

$$\begin{split} & IAB_{0|0} = p(000|000) - p(000|010) - p(100|100) - p(000|110) \\ & A \in g \quad B \in \bar{g} \implies IAB_{0|0} \leqslant 0 \end{split}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ ○ ● ●

F. J. Curchod et al., New J. Phys. 21, 023016 (2019)

CHSH inequality

$$IAB = p(00|00) - p(01|01) - p(10|10) - p(00|11) \le 0$$

Lifted term

 $IAB_{0|0} = p(000|000) - p(000|010) - p(100|100) - p(000|110)$ $A \in g \quad B \in \bar{g} \implies IAB_{0|0} \leq 0$ • Let us consider partition $g|\bar{g}$ RYSUNEK

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三 のへぐ

Centered inequalities

$$I_1^n := \sum_{i=2}^n I_{\vec{0}|\vec{0}}^{1i} - (n-2)p(\vec{0}|\vec{0}) \leqslant 0$$

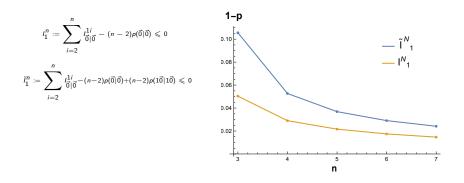
Rysunek

Symmetrical inequalities

$$I_{sym}^{n} := \sum_{i} \sum_{j>i} I_{\vec{0}|\vec{0}}^{A_{i}A_{j}} - \binom{n-1}{2} p(\vec{0}|\vec{0}) \leqslant 0$$

▲□▶ ▲□▶ ▲ 臣▶ ▲ 臣▶ ― 臣 … のへぐ

Rysunek



While in the case of 3-qubit GHZ I_1^n is resistant to around 5% of white noise, the new inequality for 10% still detects GMNL.

More general scheme of lifting allows for derivation inequalities for qudits.