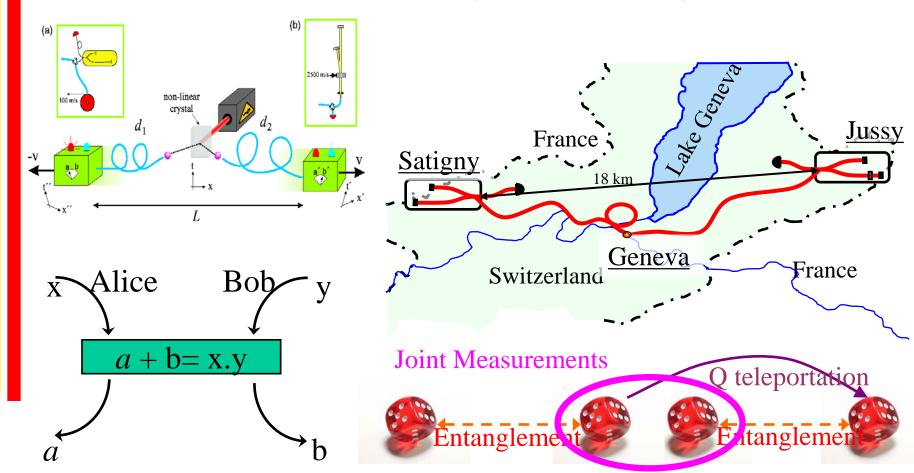


# From Bell inequalities to Quantum Measurements

#### Nicolas Gisin

Group of Applied Physics, University of Geneva Constructor University, Bremen, Germany





In the 1970's there were two main challenges affecting the foundations of quantum physics:

- 1. Quantum Non-Locality, i.e. violation of Bell inequality
- 2. The Quantum Measurement Problem, i.e. the Schrödinger eq. is interrupted whenever the system encounters a bunch of atoms with a sticker saying "measurement apparatus".

#### **General consensus:**

Quantum non-locality: kind of solved, certainly lots of progress,

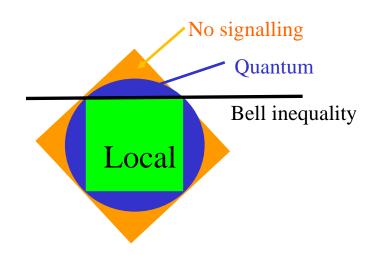
Quantum measurement problem: not much progress, certainly no agreement.



#### Bell inequalities, lots of progress

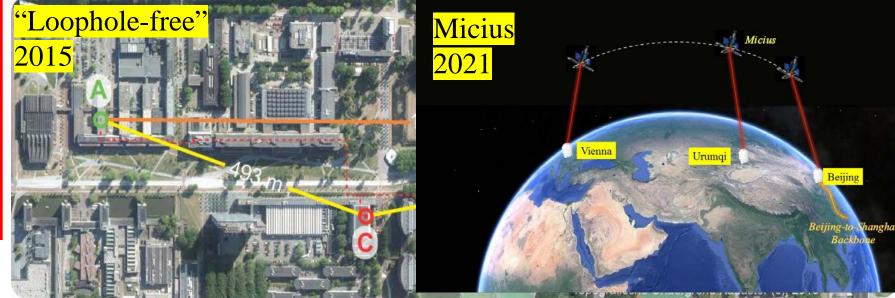
- <u>Theoretical physics</u>: communication cost to simulate quantum correlations
- Conceptual physics:
   PR-box (conceptual tool)
- Mathematical physics:
   local polytope
   ⇒ many Bell inequalities

• Experiments



#### **Experiments**



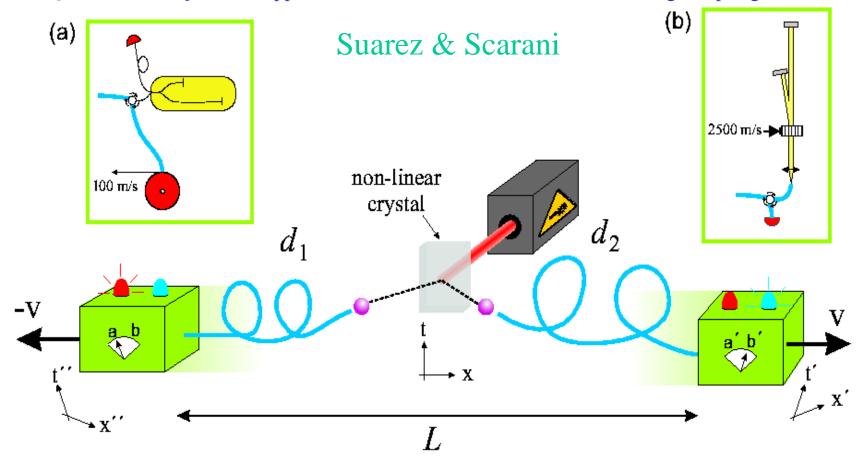




## Further experiments: before-before configurations

NG, Sundays in a quantum engineer's life, quant-ph/0104140

in Quantum [Un]speakable, pp 199-208, ed. R.A. Bertlmann and A. Zeilinger, Springer 2002



PRL <u>88</u>,120404,2002; J.Phys.A <u>34</u>,7103,2001; Phys.Lett.A <u>276</u>,1,2000



There is no c in  $\mathcal{H}$ 

Alice



Bob

Let's test whether something is going on behind the scene, as John Bell used to say.

Let's test not only common causes, but also direct causation.

All experiments so far find  $v \ge 50$  to 100 thousand times c Which goal should the community set?

c/speed sound  $\approx 10^6$ 

next generation exp should aim at  $v \approx 10^7$  c

2 orders improvement is feasible:

x10 thanks to snspd

x10 thanks to longer

distances

+ faster data acquisition





#### Bell inequalities

- Theoretic physics: communication cost to simulate quantum correlations
- Conceptual physics:
   PR-box (conceptual tool)
- Mathematical physics: local polytope ⇒ many Bell ineq
- Quantum

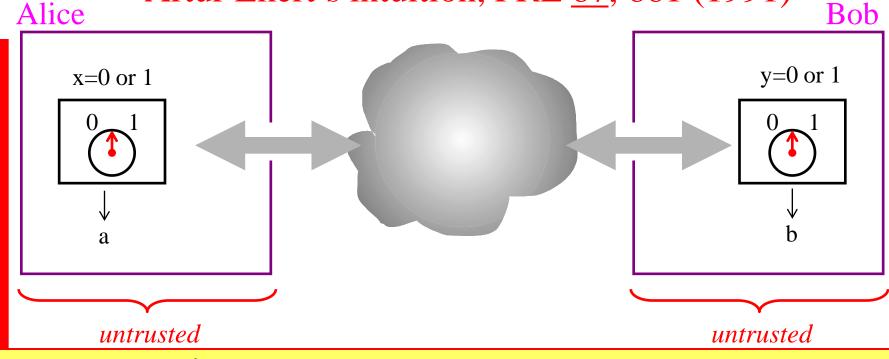
  Bell inequality

  Local

- Experiments
- Applications

#### Exploit quantum nonlocality for cryptography

Artur Ekert's intuition, PRL 67, 661 (1991)



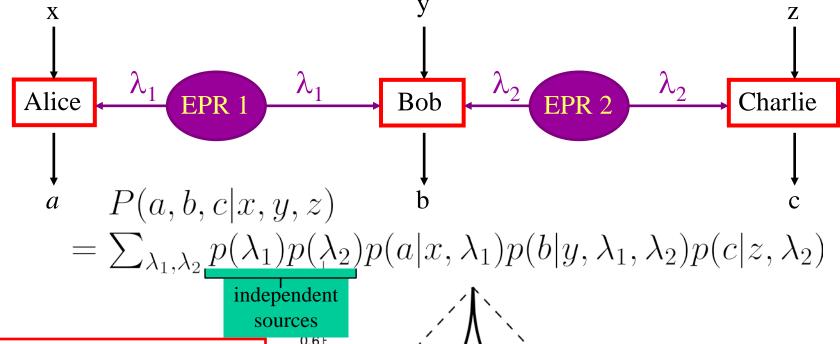
If p(a,b|x,y) violates some Bell inequality, then p(a,b|x,y) contains secrecy irrespective of any detail of the implementation!

Constructor Universit

Geneva University

#### **Network non-locality**

Branciard, NG, Pironio PRL <u>104</u>, 170401-1/4, 2010



The set of bilocal  $p_{biloc}(a,b,c \mid x,y,z)$ 

is **not convex**:



Networks with independent sources can't be described using only real Hilbert spaces!

Review: M.Renou et al., arXiv:2101.10873, Nature 6005, 625 (2021) 0.5

Tavakoli et al., Rep. Prog. Phys. 85, 056001 (2022)

 $\sqrt{|I_X| + \sqrt{|I_Y|}} \leqslant 1$ 

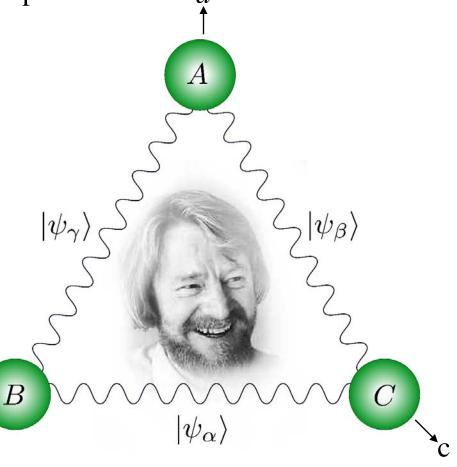


#### Triangle without inputs

P(a,b,c) satisfies the NSI principle:

- No-Signalling
- Independent sources
- 1. How to prove quantumness? Where is  $\hbar$ ?
- 2. How to prove that the outcomes contain some randomness?

Sekatski, Boreiri, Brunner, PRL 131, 100201 (2023) Boreiri et al., PRL 134, 010202 (2025)



Review on Network Non-Locality:

Tavakoli et al., Rep. Prog. Phys. 85, 056001 (2022)



## How does Nature perform the trick?

- How can these two locations out there in space-time know about each other?
- How does an event A know that it is nonlocally correlated to another event B?
- Who keeps track of who is entangled with whom?



#### Conclusion from non-locality

- ⇒ There is no spooky action at a distance: there is not a first event that influences a second event.
- ⇒ Quantum correlation just happen, without any time-ordering, somehow from outside spacetime!

(there is no story in space-time that tells us how it happens)

- ⇒ there is no mechanical explanations of quantum correlations
- ⇒ Materialism is false! (nature is not a sort of legoland)



#### The deep Q Measurement Problem

- Physics is all about extracting information about
   How Nature Does it
- Extracting information = performing measurements.

Measurements have outcomes

 Hence the Quantum Measurement Problem is a serious physics problem:

Without a resolution, Q theory is not physics.

- A resolution will lead to new physics.
- There is no reason to give up realism, nor to believe in determinism.
- realism ≠ determinism.



# Instead of fighting the deep measurement problem heads on, let's tame it, let's get more familiar with quantum measurements

Today we know a lot about

- POVMs
- Weak measurements
- Conditions under which measurements are compatible or not
- etc

But what do we know about joint measurements, measurements whose eigenstates are entangled?





# Instead of fighting the deep measurement problem heads on, let's tame it, let's get more familiar with quantum measurements

Today we know a lot about

- POVMs
- Weak measurements
- Conditions under which measurements are compatible or not
- etc

But what do we know about joint measurements, measurements whose eigenstates are entangled?

Most physicists – even within the quantum information community – would not be able to mention any joint measurement beyond the famous Bell State Measurement.

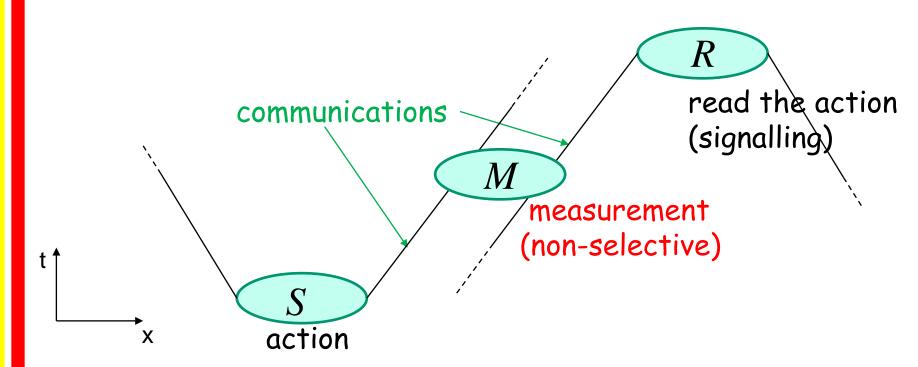


#### Impossible Measurements on Quantum Fields

RAFAEL D. SORKIN

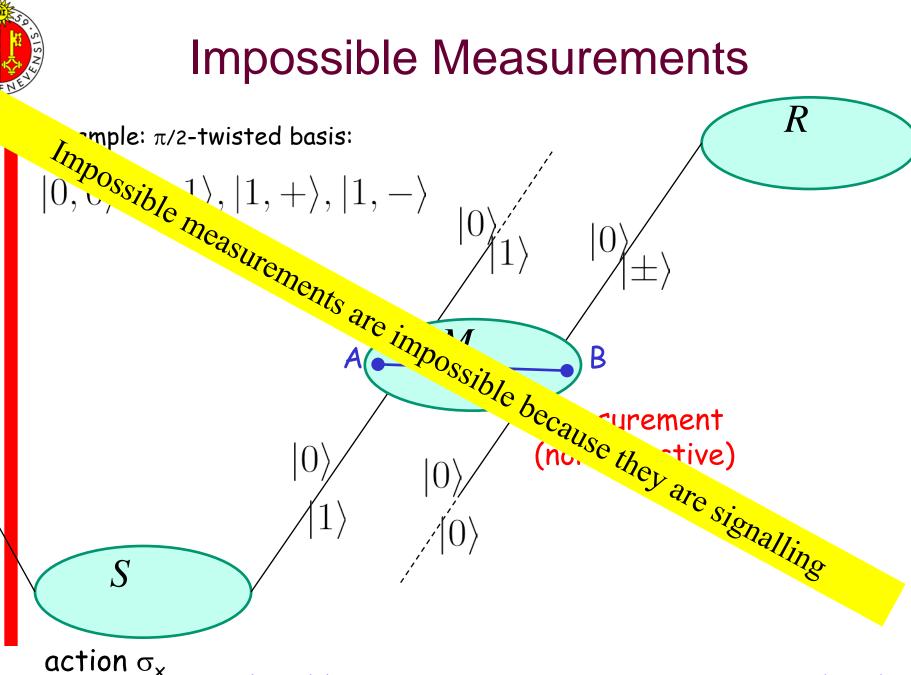
arXiv:gr-qc/9302018

Department of Physics, Syracuse University, Syracuse NY 13244-1130



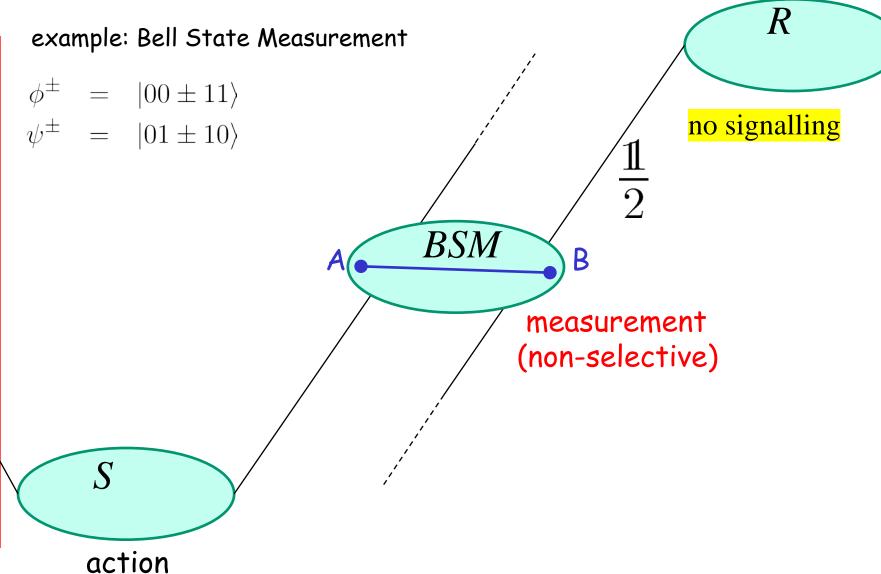
• The measurement covers some space-like region. Hence, impossible measurements are measurements of NL (non-local) variables.





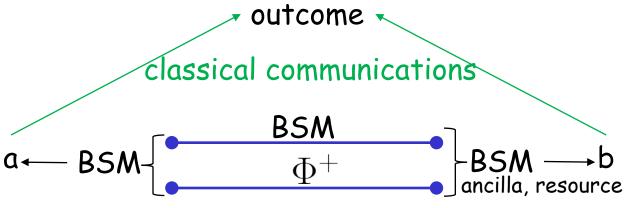


#### Impossible Measurements





#### BSM: Bell State Measurement

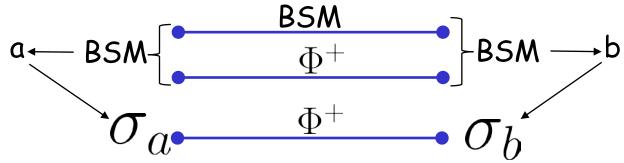


The outcomes are s.t. the number of  $\Phi$  is even and the number of + is also even.

⇒ recovers the Born rule, is not signalling, is localizable (all quantum part is done locally), but is not ideal (not immediately reproducible, not projective)



#### Ideal BSM



Now, the BSM is localized and ideal (up to some swaps).

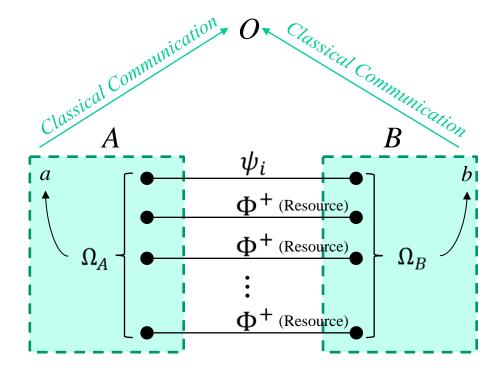
Theorem: (Popescu-Vaidman)

The BSM is the <u>only</u> joint NL measurement of 2 qubits that can be measured ideally without signalling.

Note: the BSM is not a typical measurement, but is exceptional!



#### Localizable Measurements



Theorem: All measurements, any dimension, any nb of parties, are localizable (with infinite # ebits).

(Groisman-Reznik, Vaidman)

In a nutshell: "impossible" measurements are localizable, hence possible, but can't always be ideal.

Flavio del Santo & NG, Quantum 8, 1267 (2014) Jef Pauwels, Alejandro Pozas Kerstjens, Flavio del Santo, NG, PRX 15, 021013 (2025)



#### Classification of Joint Measurements

Let's classify joint measurements in function of the entanglement cost necessary to localize them.

Obviously, with 0 entanglement only product measurements can be localized

⇒ the set of all joint quantum measurements has some structure, from the simplest to the most complex.

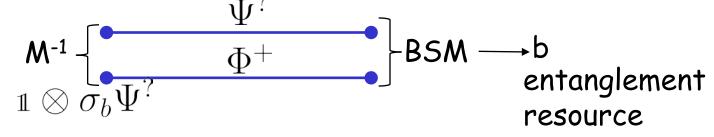
The structure is somewhat the analogue for measurements of the local polytope for quantum non-locality.

This structure is related to the Clifford hierarchy.



#### Localizable NLM with 1 ebit (2<sup>nd</sup> level)

$$M = (\psi_1 \ \psi_2 \ \psi_3 \ \psi_4)$$



$$M^{\dagger} \cdot \mathbb{1} \otimes \sigma_b \cdot M = \tilde{P}_b \Phi_b \equiv P_b$$

= permutation + phases of 00,01,10,11

Such NLM are not ideal: they are not immediately reproducible

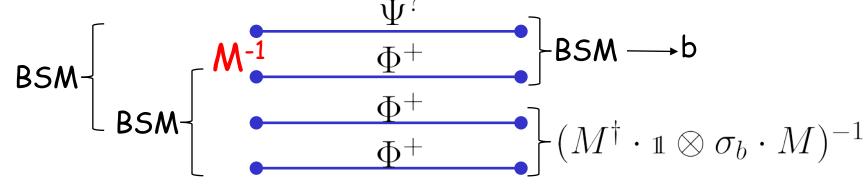
#### Theorem (2<sup>nd</sup> level of the hierarchy)

The only measurements localizable with 1 ebit, but not with 0 ebits are the

- 1.  $\pi/2$ -twisted basis measurement, and the
- Bell State Measurement.



#### Localizable NLM with 3 ebit (3<sup>rd</sup> level)



$$(M^{\dagger} \cdot \mathbb{1} \otimes \sigma_b \cdot M) \cdot \sigma_{a_1} \otimes \sigma_{a_2} \cdot (M^{\dagger} \cdot \mathbb{1} \otimes \sigma_b \cdot M)$$

= permutation+phases of 00,01,10,11

#### Theorem (3rd level of the hierarchy)

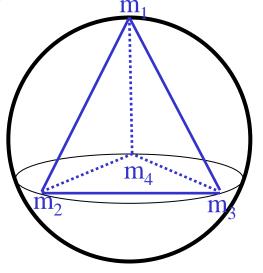
The only measurements localizable with 3 e-bit, but not with 1 e-bits are the:

- partial BSM  $\{\psi^{\pm},|00
  angle,|11
  angle\}$
- 2. elegant joint measurement (EJM), see next slide,
- another member of the EJM family arXiv:2307:06998
- 4. two members of the BS family

PR Research 6,023085(2024)



#### The Elegant Joint Measurement (EJM)



Look for 4 partially entangled and mutually orthogonal states with same degrees of entanglement and with partial states along the vertices of the tetrahedron.

$$\left| \Phi_{j} \right\rangle = c_{0} \left| \vec{m}_{j}, -\vec{m}_{j} \right\rangle + q_{0} \left| -\vec{m}_{j}, \vec{m}_{j} \right\rangle$$

$$\vec{m}_{j} = \frac{1}{\sqrt{3}} \begin{pmatrix} \frac{1}{1} & \frac{1}{-1} & -\frac{1}{1} & -\frac{1}{1} \\ \frac{1}{1} & -\frac{1}{-1} & \frac{1}{1} & -\frac{1}{1} \end{pmatrix} = \frac{\sqrt{3} + 1}{2\sqrt{2}} \left| \vec{m}_{j}, -\vec{m}_{j} \right\rangle + \frac{\sqrt{3} - 1}{2\sqrt{2}} \left| -\vec{m}_{j}, \vec{m}_{j} \right\rangle$$

$$\left\langle \Phi_{j} \middle| \Phi_{i} \right\rangle = \delta_{ji} \quad \& c_{0}, c_{1} \text{ are real}$$

$$Tr_{B}\left(\left|\Phi_{j}\right\rangle\right\langle\Phi_{j}\left|\right) = \frac{1}{2}\left(1 + \frac{\sqrt{3}}{2}\vec{m}_{j}\right) \qquad Tr_{A}\left(\left|\Phi_{j}\right\rangle\right\langle\Phi_{j}\left|\right) = \frac{1}{2}\left(1 - \frac{\sqrt{3}}{2}\vec{m}_{j}\right)$$

iso-entangled with tangle =  $\frac{1}{4}$ . Localizable, but not ideal meas.

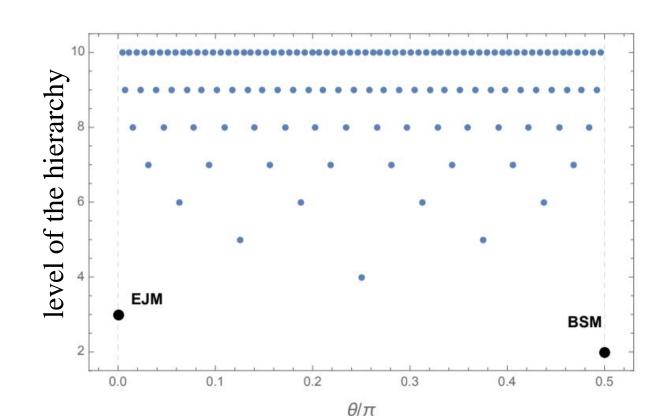


#### **Generalized EJM**

J. Pauwels, NG arXiv:250901852

J. Pauwels et al. arXiv:2509.011851

$$|\Phi_j\rangle = \frac{\sqrt{3} + e^{i\theta}}{2\sqrt{2}} |\vec{m}_j, -\vec{m}_j\rangle + \frac{\sqrt{3} - e^{i\theta}}{2\sqrt{2}} |-\vec{m}_j, \vec{m}_j\rangle$$





## The Multiqubit Elegant Joint Measurement

J. Pauwels, NG arXiv:250901852 J. Pauwels et al. arXiv:2509.011851

$$|\psi\rangle = \sum_{\mathbf{s}\in\{\pm\}^n} c_{\mathbf{s}} |s_1 \bigotimes s_2 \bigotimes \cdots \otimes s_n \bigotimes \rangle,$$

$$G_{\text{tet}}^{(n)} \equiv \left\langle Z_i Z_{i+1}, X^{\otimes n} \right\rangle \qquad \mathcal{B} = \left\{ g | \psi \rangle \,\middle|\, g \in G_{\text{tet}}^{(n)} \right\}$$

3-qubit EJM: 
$$|\psi\rangle = c_0 |-\vec{m}_1, -\vec{m}_1, -\vec{m}_1\rangle$$
  
  $+ c_1 (|\vec{m}_1, -\vec{m}_1, -\vec{m}_1\rangle + \text{perms.})$   
  $+ c_2 (|\vec{m}_1, \vec{m}_1, -\vec{m}_1\rangle + \text{perms.})$   
  $+ c_3 |\vec{m}_1, \vec{m}_1, \vec{m}_1\rangle$ .

Higher dimensions:

J. Czartowski and K. Zyczkowski, Quantum 5, 442 (2021)



#### Relation to Clifford hierarchy

J. Pauwels, NG arXiv:250901852 J. Pauwels et al. arXiv:2509.011851

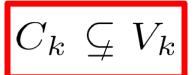
**Recursive algebraic hierarchy** organizing all unitaries by their action on Paulis

$$C_k = \{ U \in U(2^n) | UPU^{\dagger} \in C_{k-1} \ \forall P \in P_n \}, \ C_1 = P_n$$

Can be understood as complexity of nonlocal computation

Similarly, localizability is an algebraic hierarchy:

$$V_1 \subsetneq V_2 \subsetneq \ldots \subsetneq V_k \subsetneq \ldots$$



Intuition: Measurements are a subset of all computations



# What does that teach us about the deep quantum measurement problem?

May be as much as

- local polytopes,
- communication cost of simulating non-local correlations.

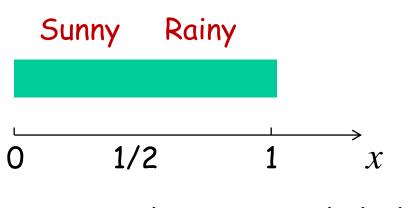
It doesn't remove the problem, but we start taming it.

#### My approach:

- is the measurement problem really new?
- aren't we asking too much to quantum theory?



#### Deterministic chaos?



 $X binary = 0.b_1b_2b_3b_4....b_n....$ 

 $\rightarrow$  X binary = 0.b<sub>2</sub>b<sub>3</sub>b<sub>4</sub>....b<sub>n</sub>.....

Whether x lies in the left half or the right half after n steps depends on the  $n^{th}$  bit,  $b_n$ , of the initial condition x.

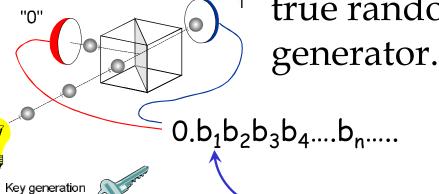
Is the billionth bits "physically real"? The question is not whether the billionth bit can be measured, but whether it corresponds to something physical?



#### Typical real numbers

- The bits of typical real numbers have no struction. When we say:

  quantial  $x_0 \in \Re$
- One rewhat we effectively say is: oall questi let  $x_0$  denote an infinite an langua amount of information.
- The only good way of thinking of a typical real number is the unlimited string of outcomes of a
  "0" True random number



G. Chaitin, The Labyrinth of the Continuum, in Meta Math!, Vintage 2008

→ typical real number



#### My intuition of Real Number

0.3906342870591187385492127064927337249582389510 87306019819422629712856540790996892145044904136<sub>2</sub> 

## TO THE PERSON OF THE PERSON OF

## Classical mathematics Real Number

0.3906342870591187385492127064927337249582389510
846240185997352602438088375931280495553838500
253876418066556042229443891573679220895949251
Instead of random events happening as time passes, 2
all the randomness is coded in the initial condition.
7
6922985461713919713178282345782969482862331709
Instead of God playing dice as time passes, God played all dies at the big-bang and coded all outcomes in the i.c.

Seems it's the mathematical language that forces us to speak of deterministic chaos!

7236099171936642858390197572002307537290953804 45847567821641758542905395011953079755477916723 19340106851299062095688236614693906421190207033 712264516218230979428454131121504352853313616564 873060198194226297128565407909968921450449041363 74815159345054159983821378228976969913337362783



Whether Newtonian classical mechanics is deterministic or not, is not a scientific question; it depends on the physical significance one associates with mathematical real numbers, i.e. it depends on the mathematical language: classical mathematics versus intuitionistic mathematics

C. Posy, Mathematical Intuitionism, Cambridge Univ. Press, 2020.



#### Intuitionistic mathematics

# Mathematical real numbers are not Physically real

### real numbers are not really real

Climate physics uses truncated numbers and stochastic remainders.

Palmer, T. N. *Nat. Rev. Phys.* **1**, 463–471 (2019). (see also his Rational Q Mechanics & implications for Q computation)

Statements about the future can be neither true nor false, but indeterminate.
 The law of the excluded middle does not hold.

The mathematical language we speak has a huge influence on the world-view that physics presents to us.

C. Posy, Mathematical Intuitionism, Cambridge Univ. Press, 2020.



#### The Classical measurement problem

- When do new bits in the series 0.b<sub>1</sub>b<sub>2</sub>b<sub>3</sub>b<sub>4</sub>...b<sub>n</sub>... gain determined values?
- Precisely, when the classical bits in the given series  $0.b_1b_2b_3b_4...b_n...$  gain relevance.
- More precisely, it is not the value of individual bits that gain determinate values, but information that gets spontaneously created, i.e. new structure of the world becomes real (true becoming).
- This is perfectly in line with the Q info slogan: information is physical.



#### The Classical hidden variables

- Postulating that the initial conditions are faithfully described by classical mathematical real numbers is an elegant way of adding all future results, while making sure that they remain inaccessible for long enough a time.
- The real numbers are the hidden (truly inaccessible) variables of classical mechanics!
- The hidden variables of classical mechanics are non-contextual, no  $\hbar$  .



## Real numbers and Bohmian mechanics

The fact is that almost all physicists accept real numbers—without noticing that they are hidden variables—while simultaneously reject Bohmian positions as unnecessary.

Instead of believing in hidden variables, let's accept indeterminacy.

As time passes, new information comes to the world and reduces indeterminacy. This leads to indeterminism and to becoming.

This solution to the classical measurement problem suggests an analogue solution to the deep quantum measurement problem:

Information is physical and this physical information grows with time, though only quantum has incompatible physical quantities, i.e.  $\hbar$ 

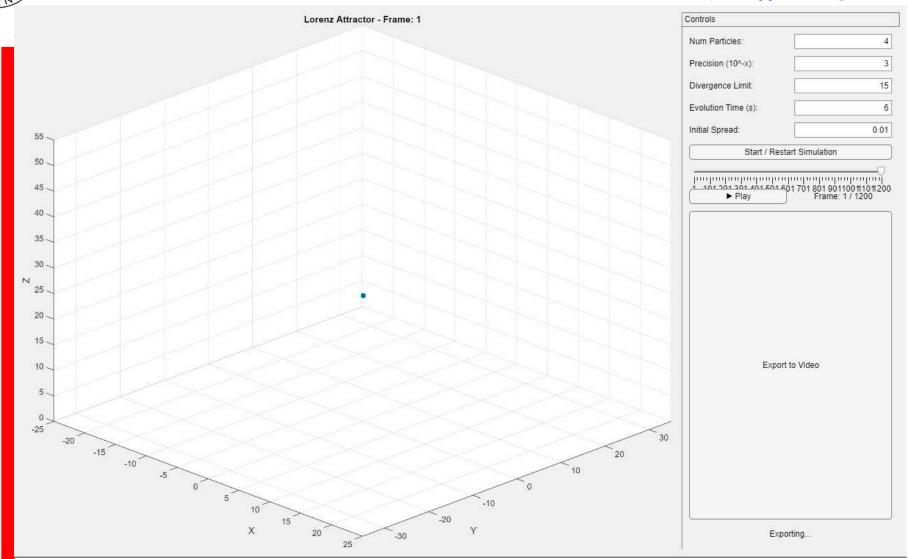
N. Gisin, *Real Numbers are the Hidden Variables of Classical Mechanics*, Quantum Studies: Mathematics and Foundations 7, 197-201 (2020). Flavio del Santo & NG, *Which features of quantum physics are not fundamentally quantum but are due to indeterminism?*, Quantum 8, 1267 (2025);

Creative and geometric times, arXiv:2404.06566



### Lorenz attractor: a classical beam-spitter

Created by Sajjad Bagdadi





#### Conclusions

All physics theories require an interpretation.

- Physical quantities are never determined with infinite precision: no infinite physical information.
- "Impossible" measurements are localizable, hence possible, but not ideally (not immediately reproducible).
- All joint quantum measurements can be classified from the simplest to the most complex by the number of e-bits necessary to localize them.
- Information is physical and this physical information grows with time.
- This applies also behind the moon: realism ≠ determinism
- Does that solve the deep measurement problem?
- No: a full solution will lead to new physics.
- There is no mechanical solution to the measurement problem